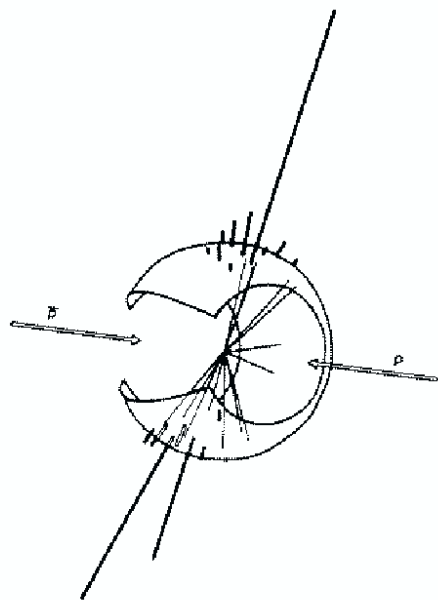


# The jet fragmentation in pp and AuAu collisions at $\sqrt{s_{NN}}=200$ GeV

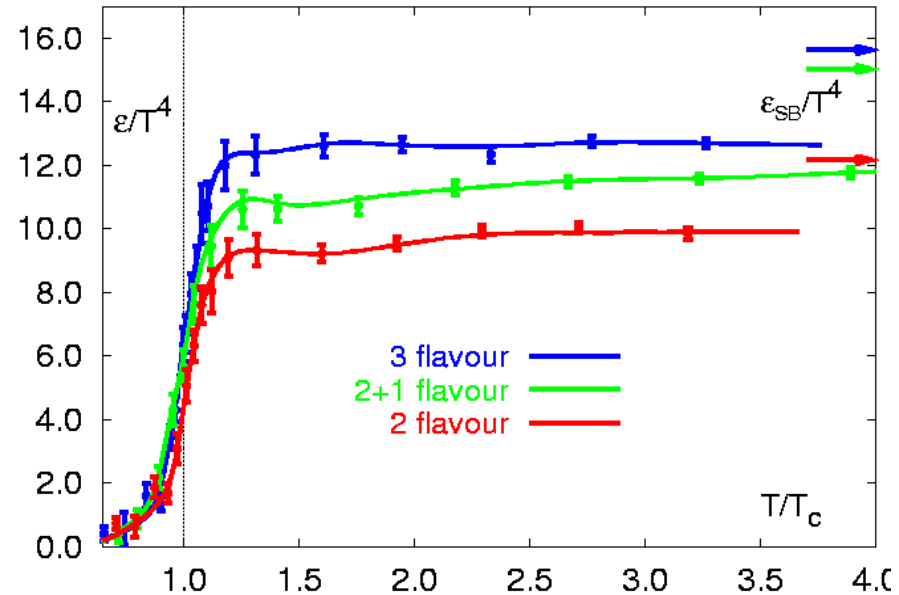


Jan Rak

University of New Mexico  
For the PHENIX collaboration

## Main RHIC results from AuAu data (with us since SB QM01 )

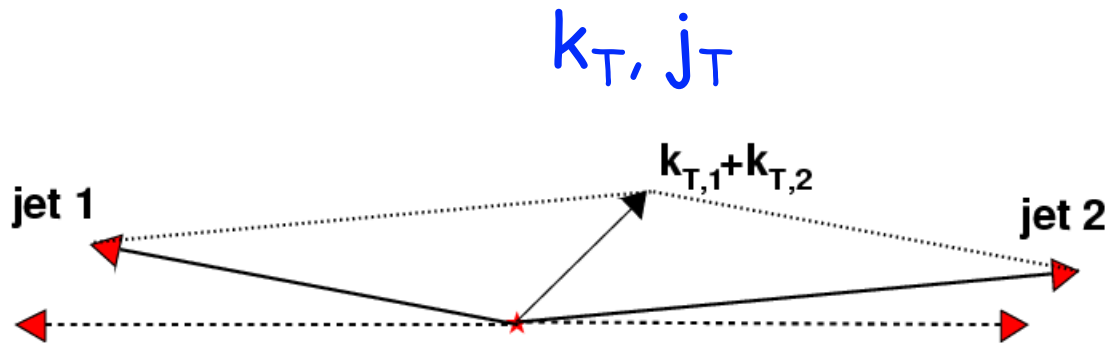
- Strong high- $p_T$  yield suppression.
- Back-to-back jet disappearance.
- Large azimuthal anisotropies.



Lattice QCD, *Lect. Notes Phys* **583**, 209 (2002)

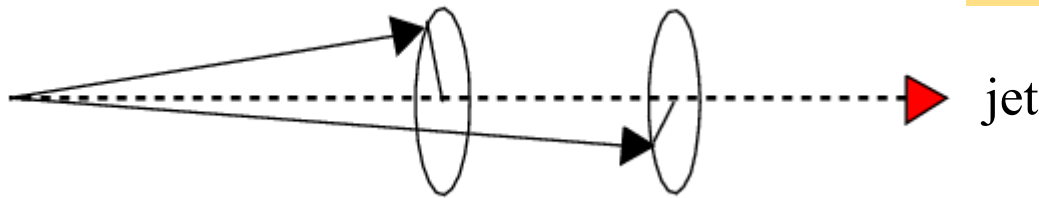
## Modification of parton properties due to their interactions with QCD medium

- parton “intrinsic” transverse momentum  $k_T$
- Jet shape – jet transverse fragmentation momentum  $j_T$
- Parton distribution function
- Fragmentation function



Partons have to materialize  
(fragment) in colorless world

$\vec{j}_T$  = jet fragmentation  
transverse momentum

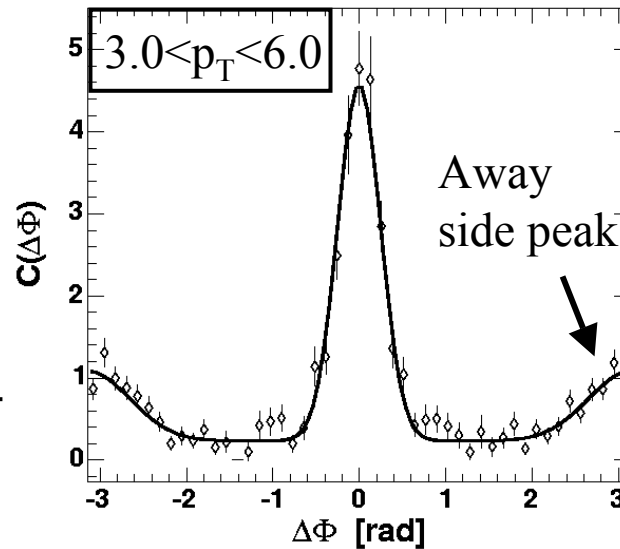
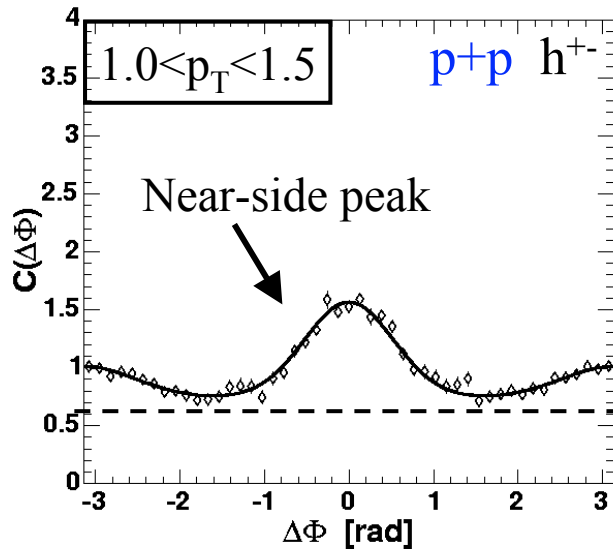


$\langle |j_{Ty}| \rangle$  is an important jet parameter. It's constant value independent on fragment's  $p_T$  is characteristic of jet fragmentation ( $j_T$ -scaling).

$\langle |k_{Ty}| \rangle$  (intrinsic + NLO radiative corrections) carries the information on the parton interaction with QCD medium.

$$\langle k_{\perp}^2 \rangle_{AA} = \langle k_{\perp}^2 \rangle_{\text{vac}} + \langle k_{\perp}^2 \rangle_{\text{IS nucl}} + \langle k_{\perp}^2 \rangle_{\text{FS nucl}}$$

# pp and dAu correlation functions

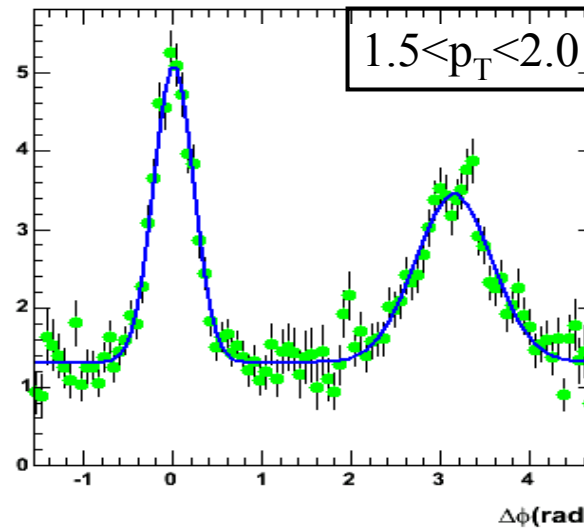
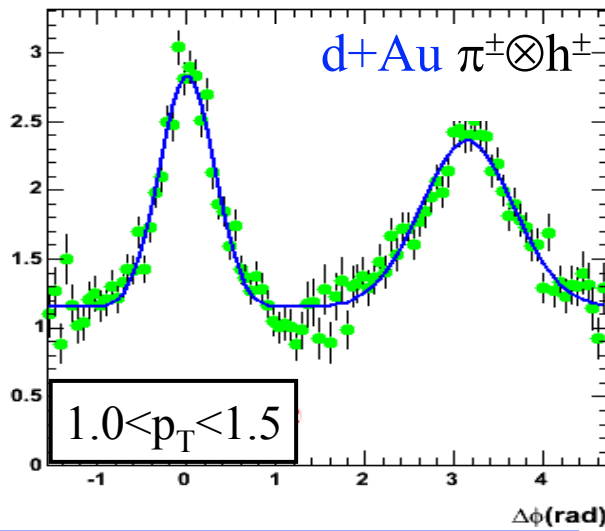


Fixed correlation:

both  $p_{T\text{trigg}}$  and  $p_{T\text{assoc}}$  are in the same range

Assorted correlation:

$p_{T\text{trigg}}$  and  $p_{T\text{assoc}}$  different



5.0 <  $p_{T\text{trigg}}$  < 16.0 GeV/c

Jet function assumed to be Gaussian

$$C_{ij}(\Delta\phi) = \text{norm} \cdot \frac{dN_{ij}^{\text{real}}}{d\Delta\phi_{ij}} / \frac{dN_{ij}^{\text{mixed}}}{d\Delta\phi_{ij}}$$



$$\text{Fit} = \text{const} + \text{Gauss}(0) + \text{Gauss}(\pm\pi)$$

## $\sigma_N, \sigma_A, \langle |j_{Ty}| \rangle, \langle |k_{Ty}| \rangle$ relations

Knowing  $\sigma_N$  and  $\sigma_A$  it is straightforward to extract  $\langle |j_{Ty}| \rangle$  and  $\langle z_{trigg} \rangle \langle |k_{Ty}| \rangle$

In the high- $p_T$  limit ( $p_T \gg \langle |j_{Ty}| \rangle$  and  $p_T \gg \langle |k_{Ty}| \rangle$ )

$$\langle |j_{\perp y}| \rangle = \langle p_{\perp} \rangle \sin \frac{\sigma_N}{\sqrt{\pi}}$$

$$\langle |k_{Ty}| \rangle \approx \langle p_T \rangle \sqrt{\sigma_A^2 - \sigma_N^2}$$

However, inspired by Feynman, Field, Fox and Tannenbaum (see *Phys. Lett. 97B (1980) 163*) we derived more accurate equation

$$\langle z_{trigg} \rangle \langle |k_{Ty}| \rangle = \frac{1}{\sqrt{2} x_h} \sqrt{\langle p_T \rangle^2 \sin^2 \sqrt{\frac{2}{\pi}} \sigma_A - (1 + \langle x_h \rangle^2) \langle |j_T| \rangle^2}$$

$$x_h = p_{T,assoc} / p_{T,trigg}$$

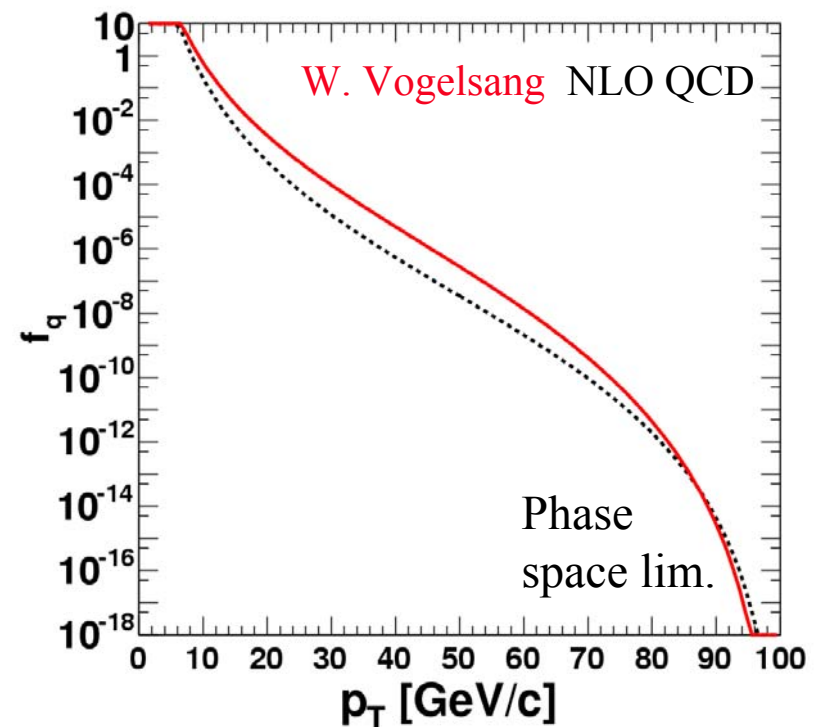
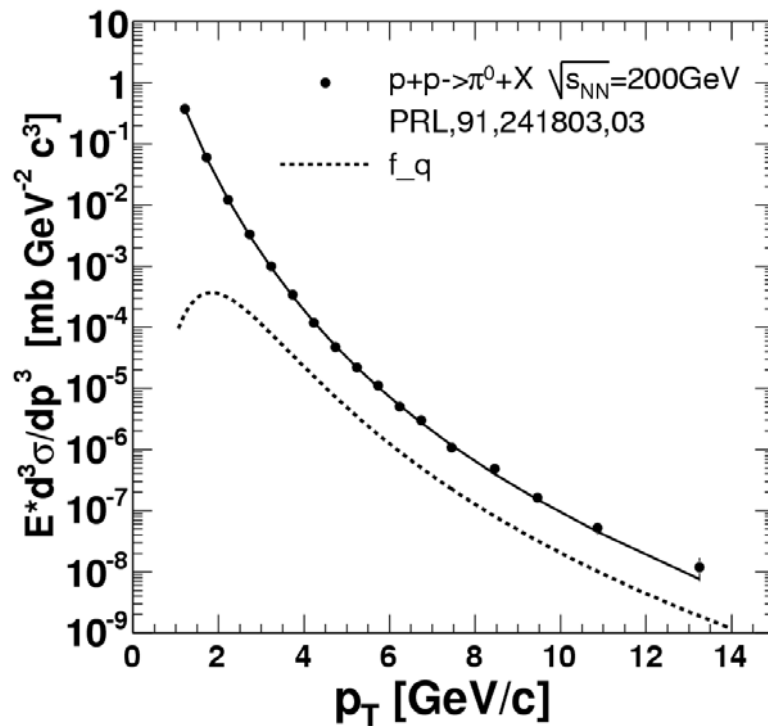
# "Effective" parton distribution function in pp

$$f_q(p_{Tq}) \propto \frac{(1 - (2p_{Tq}/\sqrt{s})^2)^8}{(M_0^2 + p_{Tq}^2)^{n/2}}$$

$$M_0 = \frac{m_0}{\ln(p_{Tq}/\lambda)}$$

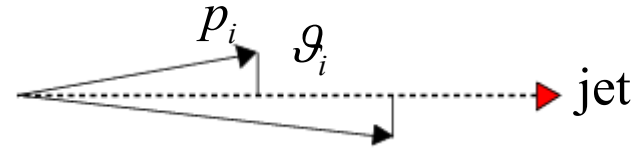
Final state parton distrib. fcn =  $\sum d\sigma_i/dt$   
 $qq \rightarrow q\bar{q}$   
 $qg \rightarrow qg$   
 $q\bar{q} \rightarrow q\bar{q}$

$$\frac{d\sigma^{\pi^0}}{p_T dp_T} = \int_{x_{Trigg}}^1 D(z) \cdot f_q(p_T / z) \cdot z^{-2} dz$$



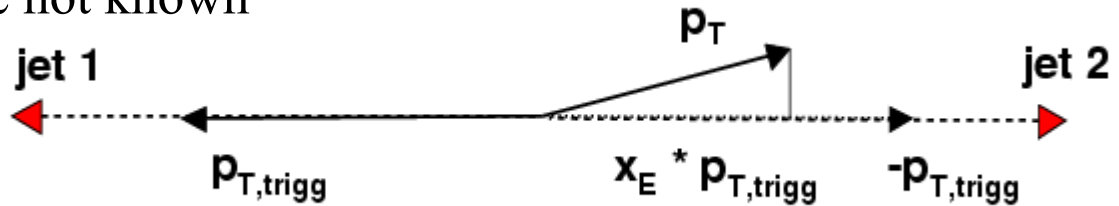
# Fragmentation Function from $x_E$ distribution

$$z_i = \frac{|\vec{p}_i| \cos(\vartheta_i)}{|\vec{p}_{parton}|} \quad \sum_i z_i = 1$$



In Practice parton momenta are not known

$$x_E = - \frac{\vec{p}_T \cdot \vec{p}_{Ttrigg}}{|\vec{p}_{Ttrigg}|^2}$$



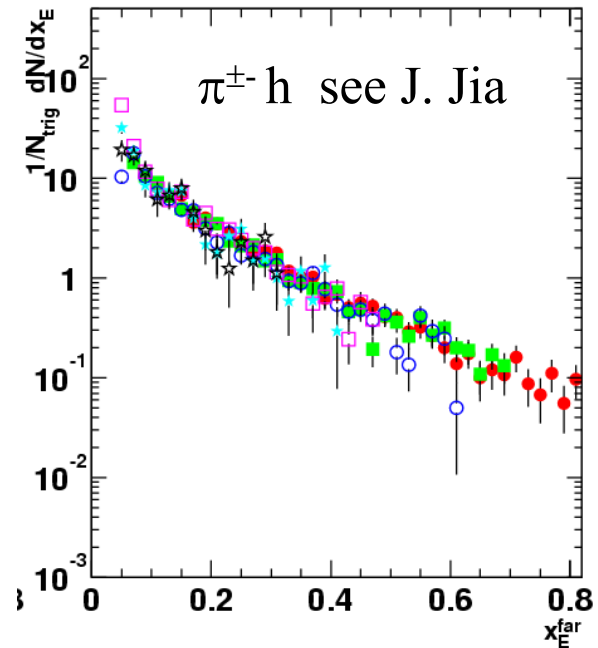
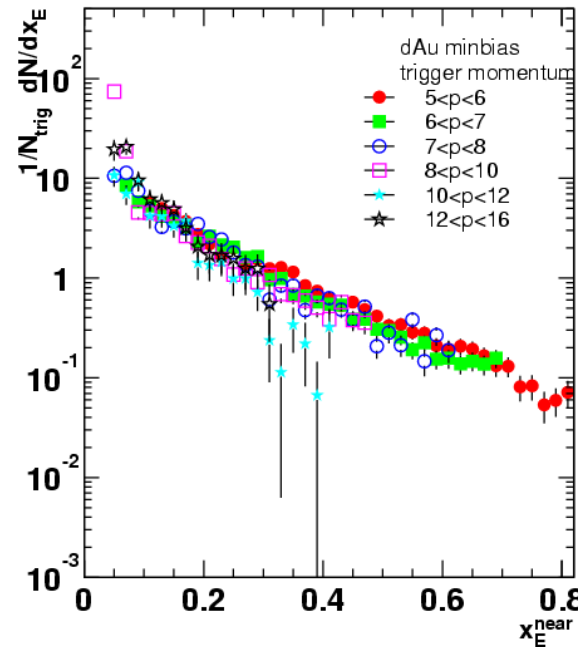
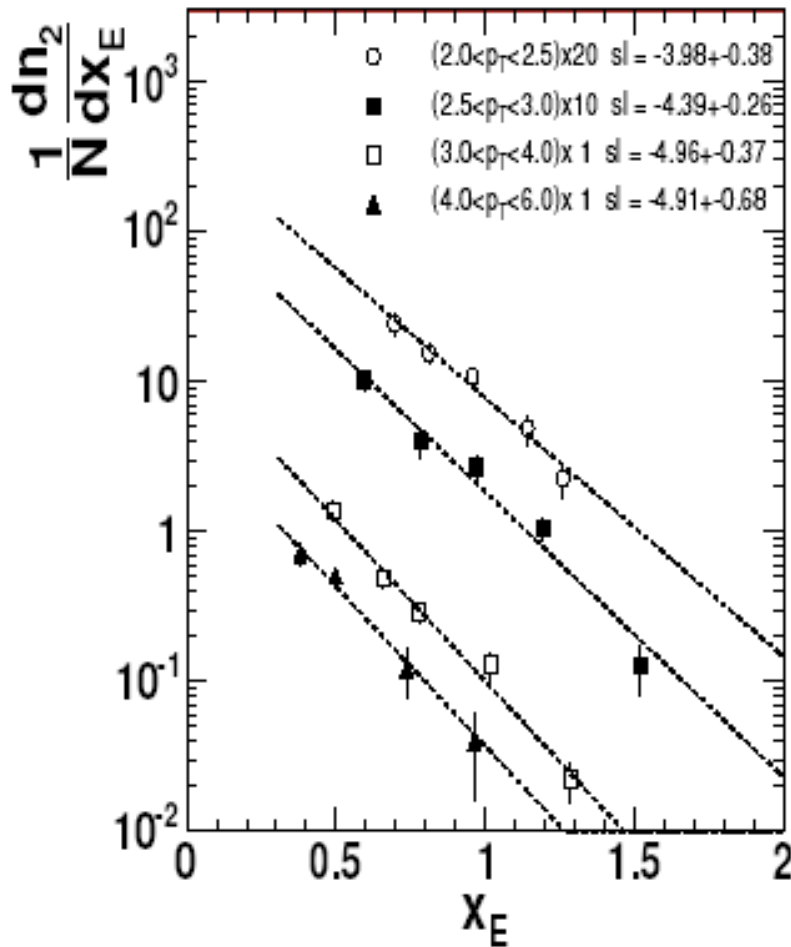
$$x_E z_{trigg} = \frac{p_T \cos(\Delta\varphi)}{p_{parton}} = z$$

⇒ Simple relation

$$\langle z \rangle = \langle x_E \rangle \langle z_{trigg} \rangle$$

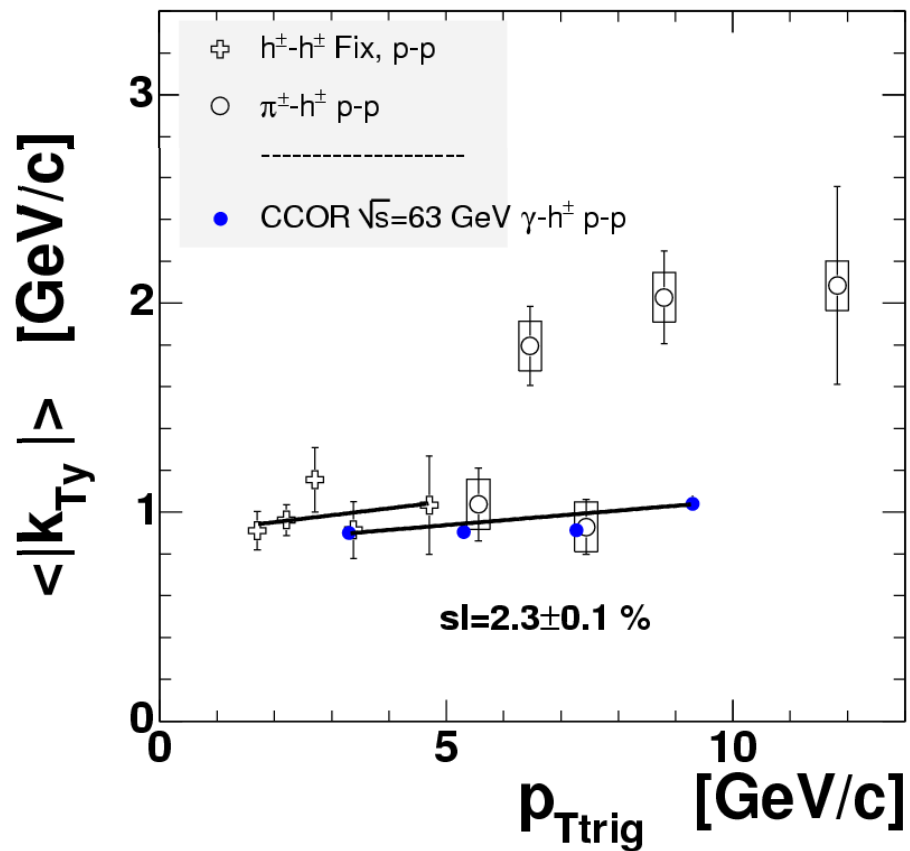
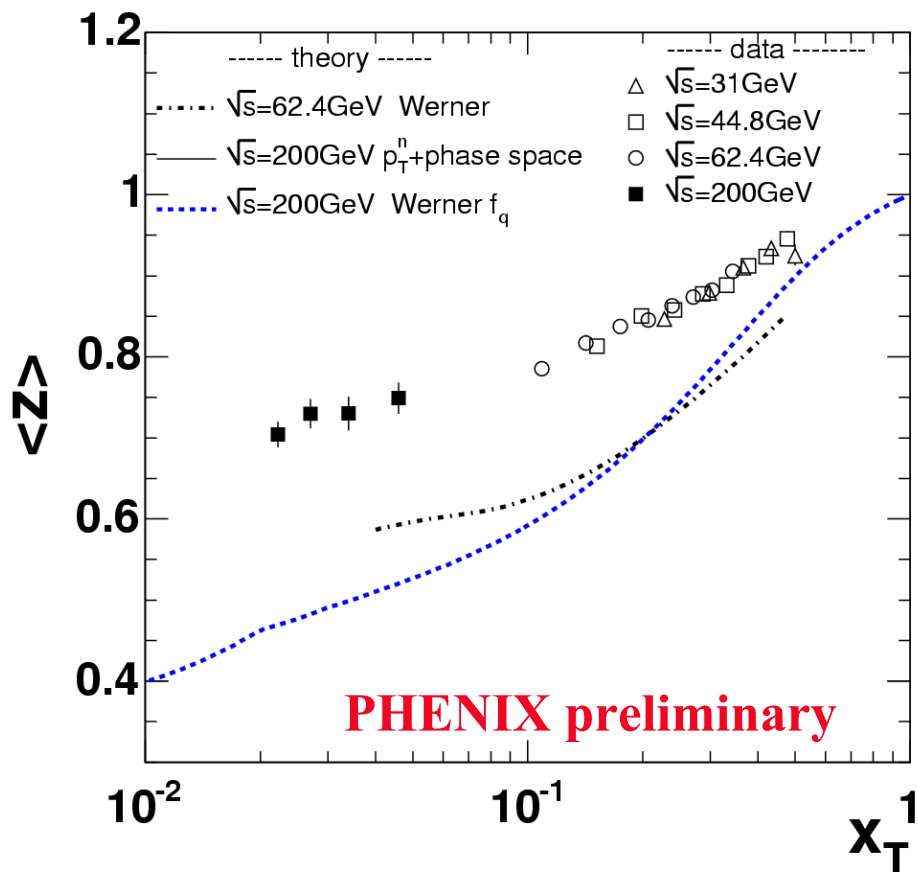
# $x_E$ in pp collisions

PHENIX preliminary





# $\langle z \rangle$ correction

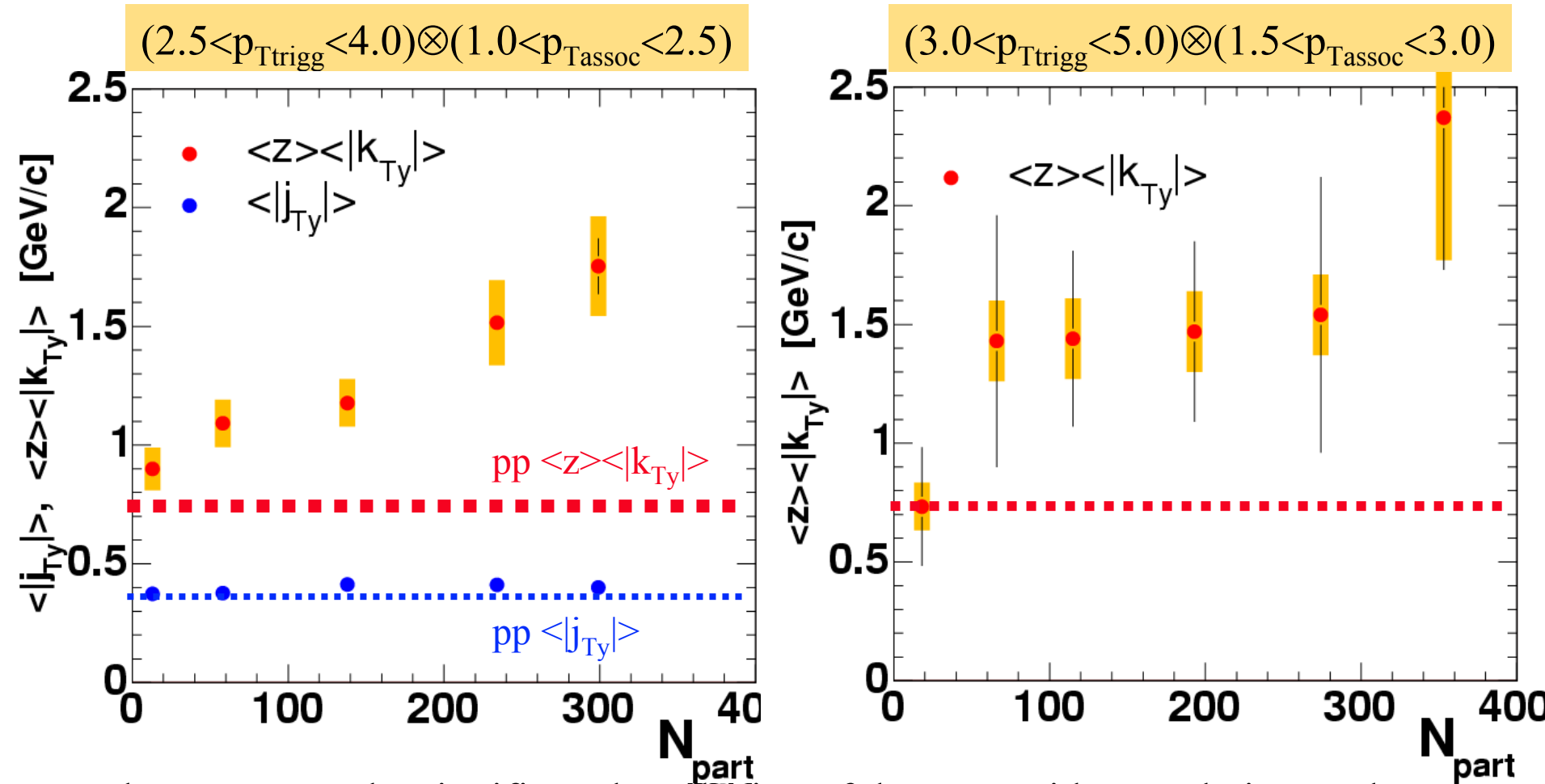


# AuAu $\langle |j_{Ty}| \rangle$ and $\langle z \rangle \langle |k_{Ty}| \rangle$ from CF

Jana Bielcikova

Phys.Rev.Lett.92:032301,2004

$$\langle k_{\perp}^2 \rangle_{AA} = \langle k_{\perp}^2 \rangle_{vac} + \langle k_{\perp}^2 \rangle_{IS\ nucl} + \langle k_{\perp}^2 \rangle_{FS\ nucl}$$



There seems to be significant broadening of the away-side correlation peak which persists also at somewhat higher  $p_T$  range.

# Summary and conclusions

## Jet production and fragmentation:

- Good agreement of the jet properties in pp collisions with other lower  $\sqrt{s}$  experiments
- dAu  $j_T$  and  $k_T$  consistent with pp
- In AuAu significant broadening of “effective”  $k_T$  - with centrality
- The pp reference understood, the analysis of  $280 \mu\text{b}^{-1}$  AuAu run-04 data is ongoing.